

Matthew Engel
 CHE 528
 TEST 2

1) Consider a binary mixture of 2 species where species 1 contains N_1 particles and species 2 contains N_2 particles. Given the potential energy, Φ , derive a formula for the second virial coefficient $B_2(T, x)$.

$$\Phi = \sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} u_{11}(r_{aa'}) + \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} u_{22}(r_{bb'}) + \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} u_{12}(r_{ab})$$

The ideal gas equation approximates atomic interactions in a dilute system. At low densities, particles are so far apart that their interactions are negligible. To mimic increasing density, or lower temperature states, a virial expansion is used to provide the first correction to an ideal gas equation of state by incorporating intermolecular interactions.

$$\frac{pV}{(N_1 + N_2)k_B T} = 1 + \frac{B_2(T, x)}{V}$$

In this case we neglect higher order terms in order to solve for the second virial coefficient.

$$p = k_B T \frac{\partial}{\partial V} \ln \int_0^\infty d\tilde{r} e^{-\beta\Phi}$$

To solve for the configurational integral, $\int_0^\infty d\tilde{r} e^{-\beta\Phi}$, one must employ the Mayer function. Where the Mayer function, $f_{ij} + 1$ can be substituted for $e^{-\beta\Phi}$.

$$e^{-\beta U(r_{ij})} = f_{ij}(r) + 1$$

Therefore, for a particular system with species 1 and 2, there are 3 possible modes of interaction in which only particles from substance 1 interact with each other (aa'), only particles from substance 2 interact with each other (bb'), and particles from species 1 and 2 interact with each other (ab).

Let $f(r_{ij}) = f_{ij}$.

$$\begin{aligned} f_{aa'} &= e^{-\beta u_{11}(r_{aa'})} - 1 \\ f_{bb'} &= e^{-\beta u_{22}(r_{bb'})} - 1 \\ f_{ab} &= e^{-\beta u_{12}(r_{ab})} - 1 \end{aligned}$$

In general, for a low density mixture the configuration integral, Q , may be written

$$\begin{aligned}
Q(N, V, T) &= \int_V dr^N e^{-\beta\Phi_N(r^N)} \\
&= \int_V dr^N \prod_{i<j}^N [1 + f(r_{ij})] \\
&= \int_V dr^N \left[1 + \sum_{i<j} f(r_{ij}) + \dots \right]
\end{aligned}$$

for each of our Mayer functions which are summed over all particles in the species. So, for a pairwise additive potential, Φ

$$\begin{aligned}
Q(N, V, T) &= \int_0^\infty dr^{\vec{2}} \prod_i \prod_{i<j} [1 + f_{ij}] \\
&= \int_0^\infty dr^{\vec{2}} \left(1 + \sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} f_{aa'} \right) \left(1 + \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} f_{bb'} \right) \left(1 + \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} f_{ab} \right)
\end{aligned}$$

Considering two unique species, where N is the number of species, we perform the integral over each species explicitly. After distributing the Mayer function the configuration integral becomes

$$\begin{aligned}
Q(N, V, T) &= \int_0^\infty dr_1^{\vec{1}} \int_0^\infty dr_2^{\vec{2}} \left(1 + \sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} f_{aa'} + \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} f_{bb'} + \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} f_{ab} \right) + \dots \\
&\dots + \left(\sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} f_{aa'} \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} f_{bb'} \right) + \dots \\
&\dots + \left(\sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} f_{aa'} \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} f_{ab} \right) + \dots \\
&\dots + \left(\sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} f_{bb'} \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} f_{ab} \right) + \dots \\
&\dots + \left(\sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} f_{aa'} \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} f_{bb'} \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} f_{ab} \right)
\end{aligned}$$

Distributing the integrals $d\vec{r}_1$ and $d\vec{r}_2$ inside the sums, the configuration integral for two species becomes

$$\begin{aligned}
Q(N, V, T) = & \left(\int_0^\infty d\vec{r}_1 \int_0^\infty d\vec{r}_2 + \sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} \int_0^\infty d\vec{r}_1 f_{aa'} + \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} \int_0^\infty d\vec{r}_2 f_{bb'} \right) + \dots \\
& \dots + \left(\sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} \int_0^\infty d\vec{r}_1 \int_0^\infty d\vec{r}_2 f_{ab} \right) + \dots \\
& \dots + \left(\sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} \int_0^\infty d\vec{r}_1 f_{aa'} \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} \int_0^\infty d\vec{r}_2 f_{bb'} \right) + \dots \\
& \dots + \left(\sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} \int_0^\infty d\vec{r}_1 f_{aa'} \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} \int_0^\infty d\vec{r}_1 \int_0^\infty d\vec{r}_2 f_{ab} \right) + \dots \\
& \dots + \left(\sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} \int_0^\infty d\vec{r}_2 f_{bb'} \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} \int_0^\infty d\vec{r}_1 \int_0^\infty d\vec{r}_2 f_{ab} \right) + \dots \\
& \dots + \left(\sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} \int_0^\infty d\vec{r}_1 f_{aa'} \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} \int_0^\infty d\vec{r}_2 f_{bb'} \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} \int_0^\infty d\vec{r}_1 \int_0^\infty d\vec{r}_2 f_{ab} \right)
\end{aligned}$$

where the volume of each species is

$$\begin{aligned}
\int_0^\infty d\vec{r}_1 &= \int d^3 r_a \int d^3 r_{a'} \dots \int d^3 r_N = V_1^N \\
\int_0^\infty d\vec{r}_2 &= \int d^3 r_b \int d^3 r_{b'} \dots \int d^3 r_N = V_2^N
\end{aligned}$$

Now recall the equation of state

$$p = k_B T \frac{\partial}{\partial V} \ln Q(N, V, T)$$

Plugging in $Q(N, V, T)$ for the configuration integral $\int_0^\infty d\vec{r} e^{-\beta\Phi}$

$$\begin{aligned}
p &= k_B T \frac{\partial}{\partial V} \ln \left(V_1^N V_2^N + \sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} \int_0^\infty d\vec{r}_1 f_{aa'} + \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} \int_0^\infty d\vec{r}_2 f_{bb'} + \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} \int_0^\infty d\vec{r}_1 \int_0^\infty d\vec{r}_2 f_{ab} \right) \\
&\dots + \left(\sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} \int_0^\infty d\vec{r}_1 f_{aa'} \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} \int_0^\infty d\vec{r}_2 f_{bb'} \right) + \dots \\
&\dots + \left(\sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} \int_0^\infty d\vec{r}_1 f_{aa'} \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} \int_0^\infty d\vec{r}_1 \int_0^\infty d\vec{r}_2 f_{ab} \right) + \dots \\
&\dots + \left(\sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} \int_0^\infty d\vec{r}_2 f_{bb'} \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} \int_0^\infty d\vec{r}_1 \int_0^\infty d\vec{r}_2 f_{ab} \right) + \dots \\
&\dots + \left(\sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} \int_0^\infty d\vec{r}_1 f_{aa'} \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} \int_0^\infty d\vec{r}_2 f_{bb'} \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} \int_0^\infty d\vec{r}_1 \int_0^\infty d\vec{r}_2 f_{ab} \right)
\end{aligned}$$

Integrating over all particles in the species

$$\begin{aligned}
p &= k_B T \frac{\partial}{\partial V} \ln \left(V_1^N V_2^N + V_1^{N_1-2} \sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} \int_0^\infty d\vec{r}_a d\vec{r}_{a'} f_{aa'} \right) + \dots \\
&\dots + \left(V_2^{N_2-2} \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} \int_0^\infty d\vec{r}_b d\vec{r}_{b'} f_{bb'} \right) + \dots \\
&\dots + \left(V_1^{N_1-2} V_2^{N_2-2} \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} \int_0^\infty d\vec{r}_a d\vec{r}_{a'} \int_0^\infty d\vec{r}_b d\vec{r}_{b'} f_{ab} \right) + \dots \\
&\dots + \left(V_1^{N_1-2} \sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} \int_0^\infty d\vec{r}_a d\vec{r}_{a'} f_{aa'} \cdot V_2^{N_2-2} \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} \int_0^\infty d\vec{r}_b d\vec{r}_{b'} f_{bb'} \right) + \dots \\
&\dots + \left(V_1^{N_1-2} \sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} \int_0^\infty d\vec{r}_a d\vec{r}_{a'} f_{aa'} \cdot V_1^{N_1-2} V_2^{N_2-2} \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} \int_0^\infty d\vec{r}_a d\vec{r}_{a'} \int_0^\infty d\vec{r}_b d\vec{r}_b f_{ab} \right) + \dots \\
&\dots + \left(V_2^{N_2-2} \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} \int_0^\infty d\vec{r}_b d\vec{r}_{b'} f_{bb'} \cdot V_1^{N_1-2} V_2^{N_2-2} \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} \int_0^\infty d\vec{r}_a d\vec{r}_{a'} \int_0^\infty d\vec{r}_b d\vec{r}_{b'} f_{ab} \right) + \dots \\
&\dots + \left(V_1^{N_1-2} \sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} \int_0^\infty d\vec{r}_a d\vec{r}_{a'} f_{aa'} \cdot V_2^{N_2-2} \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} \int_0^\infty d\vec{r}_b d\vec{r}_{b'} f_{bb'} \right) \cdot \dots \\
&\dots \cdot \left(V_1^{N_1-2} V_2^{N_2-2} \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} \int_0^\infty d\vec{r}_a d\vec{r}_{a'} \int_0^\infty d\vec{r}_b d\vec{r}_{b'} f_{ab} \right)
\end{aligned}$$

$$\begin{aligned}
p &= k_B T \frac{\partial}{\partial V} \ln \left(V_1^N V_2^N + V_1^{N_1-2} \sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} \int_0^\infty d\vec{r}_a d\vec{r}_{a'} f_{aa'} \right) + \dots \\
&\dots + \left(V_2^{N_2-2} \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} \int_0^\infty d\vec{r}_b d\vec{r}_{b'} f_{bb'} \right) + \dots \\
&\dots + \left(V_1^{N_1-2} V_2^{N_2-2} \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} \int_0^\infty d\vec{r}_a d\vec{r}_{a'} \int_0^\infty d\vec{r}_b d\vec{r}_{b'} f_{ab} \right) + \dots \\
&\dots + \left(V_1^{N_1-2} \sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} \int_0^\infty d\vec{r}_a d\vec{r}_{a'} f_{aa'} \cdot V_2^{N_2-2} \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} \int_0^\infty d\vec{r}_b d\vec{r}_{b'} f_{bb'} \right) + \dots \\
&\dots + \left(V_1^{N_1-2} \sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} \int_0^\infty d\vec{r}_a d\vec{r}_{a'} f_{aa'} \cdot V_1^{N_1-2} V_2^{N_2-2} \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} \int_0^\infty d\vec{r}_a d\vec{r}_a \int_0^\infty d\vec{r}_b d\vec{r}_b f_{ab} \right) + \dots \\
&\dots + \left(V_2^{N_2-2} \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} \int_0^\infty d\vec{r}_b d\vec{r}_{b'} f_{bb'} \cdot V_1^{N_1-2} V_2^{N_2-2} \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} \int_0^\infty d\vec{r}_a d\vec{r}_{a'} \int_0^\infty d\vec{r}_b d\vec{r}_{b'} f_{ab} \right) + \dots \\
&\dots + \left(V_1^{N_1-2} \sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} \int_0^\infty d\vec{r}_a d\vec{r}_{a'} f_{aa'} \cdot V_2^{N_2-2} \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} \int_0^\infty d\vec{r}_b d\vec{r}_{b'} f_{bb'} \right) \cdot \dots \\
&\dots \cdot \left(V_1^{N_1-2} V_2^{N_2-2} \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} \int_0^\infty d\vec{r}_a d\vec{r}_{a'} \int_0^\infty d\vec{r}_b d\vec{r}_{b'} f_{ab} \right)
\end{aligned}$$

Allow the sum of pairwise-additive interactions to be expressed in terms of number of particles. This method accounts for different volumes of species 1 and species 2 and reflects this in the following manner

$$\begin{aligned}
\Phi &= \sum_{a=1}^{N_1} \sum_{a'=a+1}^{N_1} u_{11}(r_{aa'}) + \sum_{b=N_1+1}^{N_1+N_2} \sum_{b'=b+1}^{N_1+N_2} u_{22}(r_{bb'}) + \sum_{a=1}^{N_1} \sum_{b=N_1+1}^{N_1+N_2} u_{12}(r_{ab}) \\
&= \left(\frac{N_1(N_1-1)}{2} \right) u_{11}(r_{aa'}) + \left(\frac{(N_1+N_2)(N_1+N_2-1)}{2} \right) u_{22}(r_{bb'}) + (N_1(N_2-1)) u_{12}(r_{ab})
\end{aligned}$$